

# Unique Role of Hot-Electron Induced Self-Heating in Determining Gate-Stack Dependent Dynamic $R_{ON}$ of AlGaIn/GaN HEMTs Under Semi-ON State

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**Abstract**—In this work, we report a unique dependence of dynamic  $R_{ON}$  on gate-stack of AlGaIn/GaN high electron mobility transistors (HEMTs) on carbon (C)-doped GaN buffer under semi-on condition. Unlike the OFF-state stress, semi-ON state stressing of the devices revealed a significantly lower dynamic  $R_{ON}$  in the Schottky HEMTs when compared to SiN<sub>x</sub>-gated HEMTs. Detailed experiments including substrate bias dependence, electro-luminescence (EL) spectrum, and electric field analysis established electron trapping in the GaN buffer to be the source of dynamic  $R_{ON}$  in both the HEMT types. An on-the-fly thermal analysis, using thermoreflectance technique, revealed Schottky HEMTs to have a higher hot electron-induced self-heating near the field plate (FP) edge as compared to SiN<sub>x</sub>-gated devices. The higher lattice temperature was found to result in a significantly higher electron de-trapping from the buffer traps. The higher de-trapping reduced the dynamic  $R_{ON}$  in the Schottky HEMTs as compared the SiN<sub>x</sub>-gated devices.

**Index Terms**—AlGaIn/GaN high electron mobility transistors (HEMTs), carbon (C)-doped buffer, dynamic ON-resistance, hot electrons, self-heating, semi-ON state.

## I. INTRODUCTION

WIDE scale deployability of AlGaIn/GaN high electron mobility transistors (HEMTs) for power applications is still limited, despite their superior performance. This is majorly due to the increased ON-resistance ( $R_{ON}$ ) exhibited by these devices immediately after switching from high-voltage OFF-state to ON-state [1], [2], [3], [4], [5], [6], [7], [8], [9], [10]. Significant efforts have been made in the past to understand factors affecting this increase in  $R_{ON}$ , referred to as

Manuscript received 3 August 2022; revised 19 September 2022; accepted 2 October 2022. Date of publication 19 October 2022; date of current version 30 November 2022. This work was supported by the Department of Science and Technology (DST), Government of India, and carried out at the Indian Institute of Science, Bengaluru, under Project SP/DSTO-21-0135. The review of this article was arranged by Editor M. Hua. (Corresponding author: Sayak Dutta Gupta.)

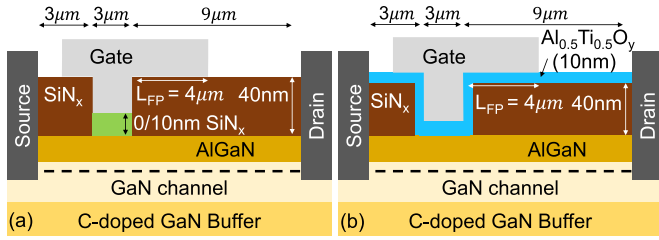
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Color versions of one or more figures in this article are available at <https://doi.org/10.1109/TED.2022.3212327>.

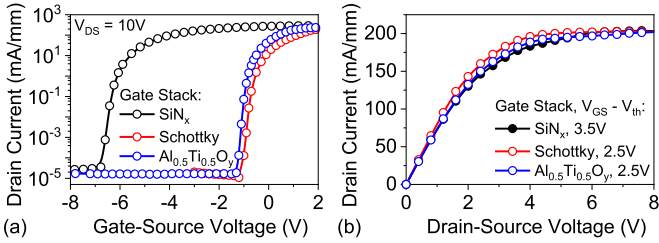
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dynamic  $R_{ON}$ , upon stressing in the OFF-state. This has led to several surface passivation approaches using SiN<sub>x</sub> [1], AlN [2], GaN cap [3], and AlTiO [4], to help suppress the dynamic  $R_{ON}$  of the devices. Besides, GaN buffer designs [7], [8], [9] have also been reported to reduce the dynamic  $R_{ON}$  of these devices. However, some recent works have highlighted that not only OFF-state stressing, but also the stress conditions applied on the devices during hard switching can significantly affect its  $R_{ON}$  [11], [12], [13], [14]. During hard switching, the device is subjected to semi-ON state conditions, wherein both high electric field and high carrier density exist simultaneously [14]. This triggers the generation of hot electrons, a major source of reliability challenges in GaN HEMTs [11], [12], [13], [14], [15], [16], [17], [18], with dynamic  $R_{ON}$  being one of the most prominent ones [11], [12], [13], [14], [15]. Indeed, several works reported a higher dynamic  $R_{ON}$  due to semi-on stressing of GaN HEMTs, as compared to OFF-state stress [11], [12], [13], [14]. Recent works have also highlighted that trapping during semi-ON state is mostly at the device surface or in AlGaIn barrier, with negligible trapping in the GaN buffer [12], [13]. However, these studies are mostly limited to lower drain to source voltages ( $V_{DS-Stress} \leq 50$  V). Our recent studies have highlighted that stressing the device at higher  $V_{DS-Stress}$  in OFF-state can lead to significant electron trapping in buffer and result in very high dynamic  $R_{ON}$  [4], [5], [6]. Moreover, dynamic  $R_{ON}$  in these stressing conditions was found to be a function of stress time ( $t_{Stress}$ ). This makes it imperative to evaluate impact of electron trapping in the GaN buffer under semi-ON state stress, as a function of  $V_{DS-Stress}$  and  $t_{Stress}$ .

In this work, we report and probe a unique  $V_{DS-Stress}$  and  $t_{Stress}$ -dependent dynamic  $R_{ON}$  behavior in AlGaIn/GaN HEMTs on carbon (C)-doped GaN buffer under semi-on conditions, which is also a function of the gate-stack. Experiments with varying substrate bias and electro-luminescence (EL) spectrum analysis are carried out to evaluate role of trapping in the GaN buffer and on the device surface, in determining dynamic  $R_{ON}$ . Furthermore, devices with relaxed electric field near the field plate (FP) edge, using Al<sub>0.5</sub>Ti<sub>0.5</sub>O<sub>y</sub> surface protection [4] are tested to determine impact of electric field



**Fig. 1.** Schematic of the AlGaIn/GaN HEMTs fabricated for this work. (a) Schottky-gated (depicted by 0 nm SiN<sub>x</sub>) and 10 nm SiN<sub>x</sub>-gated HEMTs both with 40 nm SiN<sub>x</sub> passivation. (b) Al<sub>0.5</sub>Ti<sub>0.5</sub>O<sub>y</sub>-gated HEMTs with 10 nm Al<sub>0.5</sub>Ti<sub>0.5</sub>O<sub>y</sub> deposited over 40 nm SiN<sub>x</sub> passivation.

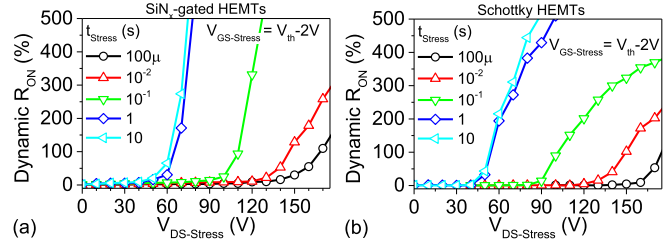


**Fig. 2.** (a) Transfer and (b) output characteristics of the fabricated AlGaIn/GaN HEMTs. A higher gate overdrive ( $V_{GS} - V_{th}$ ) was provided to the SiN<sub>x</sub>-gated HEMTs to ensure similar  $R_{ON}$  as Schottky and Al<sub>0.5</sub>Ti<sub>0.5</sub>O<sub>y</sub>-gated HEMTs, in spite of its weaker gate control.

near FP edge on trapping in GaN buffer and resulting dynamic  $R_{ON}$ . Finally, channel temperature measurements are carried out to determine any contribution of hot electron-induced self-heating in the device on its dynamic  $R_{ON}$  behavior. The electric field and hot electron-induced self-heating near the FP edge of the HEMTs, besides  $t_{Stress}$ , are established to be the factors governing the dynamic  $R_{ON}$  behavior of the HEMTs. The manuscript is organized as follows. In Section II, the device fabrication and experimental setup is described. The dynamic  $R_{ON}$  behavior of the fabricated devices is then discussed in Section III. The factors governing the dynamic  $R_{ON}$  behavior under semi-ON state stress and the related phenomenon are discussed in Sections IV and V, respectively. Finally, the work is concluded in Section VI.

## II. DEVICE FABRICATION AND EXPERIMENTATION

AlGaIn/GaN HEMTs were fabricated on a commercial grade C-doped GaN on Si epi-stack with an in-situ 40 nm SiN<sub>x</sub> passivation. The epi-stack had a 2-D electron gas (2DEG) density of  $\sim 8 \times 10^{12} \text{ cm}^{-2}$  at the AlGaIn/GaN hetero-interface. The following device variants were fabricated on the epi-stack using a well-optimized processing technology, discussed in detail in our previous works [4], [5], [19]: 1) Schottky HEMTs; 2) 10 nm SiN<sub>x</sub>-gated HEMTs, both with 40 nm SiN<sub>x</sub> passivation, as seen in Fig. 1(a); and 3) 10 nm Al<sub>0.5</sub>Ti<sub>0.5</sub>O<sub>y</sub> (AlTiO)-gated HEMTs with 40 nm SiN<sub>x</sub> passivation + 10 nm AlTiO surface protection layer (AlTiO/SiN passivation), seen in Fig. 1(b). It is worth highlighting here that the devices were processed simultaneously using a well-optimized process to minimize process-related device variabilities. Fig. 2(a) and (b) show the transfer and output characteristics of the fabricated devices. The threshold voltage ( $V_{th}$ ) of the SiN<sub>x</sub>, Schottky, and AlTiO-gated devices was  $-6.5$ ,  $-0.9$ , and  $-1.1$  V, respectively,  $V_{th}$  being measured as the  $V_{GS}$  for  $I_D = 1 \mu\text{A}/\text{mm}$ .



**Fig. 3.** Dynamic  $R_{ON}$  of (a) SiN<sub>x</sub>-gated and (b) Schottky HEMTs with  $V_{DS-Stress}$ , as a function of stress time ( $t_{Stress}$ ) under OFF-state stressing.

While the high- $\kappa$  ( $\kappa \sim 25$ ) and p-type nature of AlTiO resulted in  $V_{th}$  of  $-1.1$  V for the AlTiO-gated HEMTs [4], [19], weaker channel control due to the use of 10 nm thick low- $\kappa$  ( $\kappa \sim 5$ ) SiN<sub>x</sub> gate dielectric resulted in the significantly negative  $V_{th}$  of  $-6.5$  V for the SiN<sub>x</sub>-gated HEMTs [4]. Furthermore, proper surface cleaning before deposition of gate metal resulted in  $V_{th}$  of  $-0.9$  V for the Schottky HEMTs [5]. Moreover, the HEMTs showed good ON-state and OFF-state performance with  $I_{ON}/I_{OFF} \sim 10^7$ .

Dynamic  $R_{ON}$  behavior of the fabricated devices was evaluated by following a measure–stress–measure–relax routine [4], [5]. The HEMT's pristine  $R_{ON}$  ( $R_{Pristine}$ ) was first measured from inverse slope of  $I_D - V_{DS}$  for  $V_{DS} = 0.25 - 0.5$  V. The devices were then stressed under semi-on condition with  $V_{GS-Stress} = V_{th} + 0.3$  V, which resulted in  $\sim 3$  orders increase in device  $I_D$  as compared to OFF-state. The output characteristics of the HEMTs were then measured  $\sim 5$  ms post the stress, to evaluate its post-stress  $R_{ON}$  ( $R_{Post-Stress}$ ). Dynamic  $R_{ON}$  of the devices was then measured as: Dynamic  $R_{ON} = (R_{Post-Stress} - R_{Pristine})/R_{Pristine} \times 100$  (%). Further,  $V_{DS-Stress}$  was varied from 0 to 200 V during the stressing, with  $t_{Stress}$  ranging from 100  $\mu\text{s}$  to 20 s.

The dynamic  $R_{ON}$  of the HEMTs due to OFF-state stressing was also evaluated by biasing the gate at  $V_{GS-Stress} = V_{th} - 2$  V, to compare the behavior with the semi-ON state dynamic  $R_{ON}$ . Moreover, given that the stressing results in degradation of the device  $R_{ON}$ , the devices were allowed a 300 s relaxation period between the  $R_{Post-Stress}$  measurement and the next stress cycle to completely recover the device.

## III. DYNAMIC $R_{ON}$ BEHAVIOR OF ALGAN/GAN HEMTs

The dynamic  $R_{ON}$  behavior of SiN<sub>x</sub>- and Schottky-gated GaN HEMTs under OFF-state stress is shown in Fig. 3(a) and (b), respectively. The devices show presence of a  $t_{Stress}$ -dependent critical  $V_{DS-Stress}$  ( $V_{cr}$ ) beyond which the dynamic  $R_{ON}$  increased significantly ( $>25\%$ ). This dynamic  $R_{ON}$  behavior can be attributed to electron trapping in C-doping-induced traps in the GaN buffer during OFF-state stressing [4], [5]. Further, to determine how electron trapping affects dynamic  $R_{ON}$  behavior under semi-ON state stress, the SiN<sub>x</sub>- and Schottky-gated devices were subjected to the measure–stress–measure routine in semi-ON state bias condition to extract their dynamic  $R_{ON}$ , as seen in Fig. 4(a) and (b), respectively. The figures indicate that the devices exhibit a  $t_{Stress}$ -dependent  $V_{cr}$  under semi-ON state stressing as well. The  $V_{cr}$  initially reduces with increase in  $t_{Stress}$  and remains unchanged as  $t_{Stress}$  is increased beyond

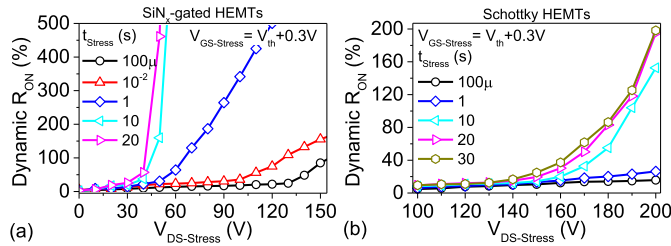


Fig. 4. Dynamic  $R_{ON}$  of (a) SiN<sub>x</sub>-gated and (b) Schottky HEMTs with  $V_{DS-Stress}$ , as a function of stress time under semi-ON state stressing.

10's of seconds. However, the SiN<sub>x</sub>-gated devices exhibit a significantly lower  $V_{cr}$  and the onset of increase in dynamic  $R_{ON}$  for lower  $t_{Stress}$  values, as compared to Schottky-gated devices. A comparison of Figs. 3 and 4 thus suggests electron trapping in C-doped GaN buffer to be a critical phenomenon determining dynamic  $R_{ON}$  behavior of the device in semi-ON state as well.

The  $t_{Stress}$  dependence of dynamic  $R_{ON}$  in SiN<sub>x</sub>- and Schottky-gated devices under OFF and semi-ON state stressing is worth highlighting here. Fig. 3(a) and (b) show that under OFF-state stressing both SiN<sub>x</sub>- and Schottky-gated devices show a similar  $t_{Stress}$  dependence with a minor variation in  $V_{cr}$ . On the other hand, Fig. 4 shows that while SiN<sub>x</sub>-gated devices exhibited a  $V_{cr}$  like phenomenon for  $t_{Stress}$  as low as 100  $\mu$ s, the Schottky HEMTs did not show any  $V_{cr}$  for  $t_{Stress} < 10$  s. This establishes the presence of an additional phenomenon in semi-ON state, which is not observed in OFF-state, that controls dynamic  $R_{ON}$  behavior of the device. Furthermore, this additional process in semi-ON state exhibits a gate-stack dependence leading to different dynamic  $R_{ON}$  behavior of Schottky- and SiN<sub>x</sub>-gated devices.

#### IV. FACTORS AFFECTING DYNAMIC $R_{ON}$ UNDER SEMI-ON STATE

Discussions in previous section suggested the presence of electron trapping-induced  $V_{cr}$  in dynamic  $R_{ON}$  behavior of GaN HEMTs under semi-ON state stressing. Moreover, the gate-stack dependence of  $t_{Stress}$  behavior of dynamic  $R_{ON}$  in semi-ON state stressing condition suggested presence of an additional phenomenon. In this section, detailed experimentations are carried out to probe into different factors affecting  $V_{cr}$  and dynamic  $R_{ON}$  of the device under semi-ON state stressing. All measurements in this section are carried out under semi-ON state stressing unless stated otherwise.

##### A. Source of Dynamic $R_{ON}$

Dynamic  $R_{ON}$  in AlGaIn/GaN HEMTs is attributed to trapping in surface and/or buffer traps [11]. To segregate the role of buffer and surface traps in determining the dynamic  $R_{ON}$  of GaN HEMTs, determination of dynamic  $R_{ON}$  behavior as a function of substrate bias ( $V_{Sub}$ ) can play a critical role [4], [5]. Modification of  $V_{Sub}$  would have negligible impact on surface trapping-induced dynamic  $R_{ON}$  due to screening by the 2DEG. However, changing  $V_{Sub}$  would significantly vary carrier injection and thereby trapping in the buffer or buffer-induced dynamic  $R_{ON}$  [4], [5]. Fig. 5(a) shows the dynamic  $R_{ON}$  of the

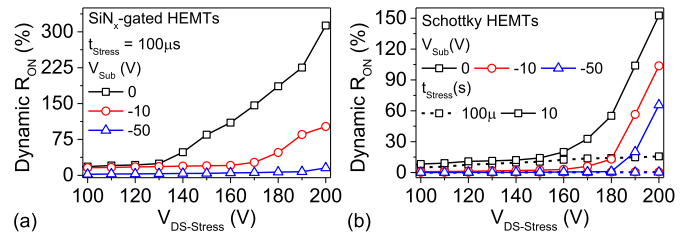


Fig. 5. Dynamic  $R_{ON}$  under semi-ON state stressing of (a) SiN<sub>x</sub>- and (b) Schottky HEMTs with  $V_{DS-Stress}$ , as a function of substrate bias ( $V_{Sub}$ ).

SiN<sub>x</sub>-gated HEMTs to significantly reduce as  $V_{Sub}$  is reduced below 0 V. A similar dependence of dynamic  $R_{ON}$  on  $V_{Sub}$  is noticed for the Schottky HEMTs in Fig. 5(b). This correlation between the HEMT's dynamic  $R_{ON}$  and  $V_{Sub}$  indicates a greater role of buffer traps in the observed dynamic  $R_{ON}$  behavior.

Trapping in the gate-stack can be another source of gate-stack dependent dynamic  $R_{ON}$  in GaN HEMTs. To determine any trapping in the gate-stack, transfer characteristics of the HEMTs was measured before and after the stress cycle. The results, as shown in Fig. 6(a), depict a negligible impact of stress cycle on  $V_{th}$  of the devices. This indicates trapping to be more prominent in the access region as compared to the gate-stack of the device.

Above discussion establishes trapping in access region of the C-doped GaN buffer to be the major contributor in determining the HEMT's dynamic  $R_{ON}$  behavior. The mechanism responsible for the trapping is discussed in the next section.

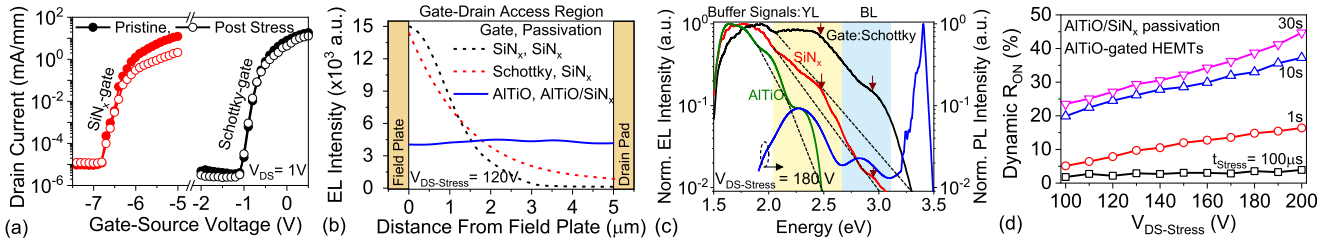
##### B. Mechanism Responsible for Trapping

Given the role of buffer traps in the observed dynamic  $R_{ON}$  behavior, there can be two major mechanisms governing the phenomenon: 1) electron trapping [4], [5] and 2) hole emission and redistribution [10]. Electron trapping refers to the phenomenon where channel electrons injected into the C-doped GaN buffer get trapped in the buffer acceptor traps, thereby resulting in dynamic  $R_{ON}$  [5]. On the other hand, emission of holes from buffer acceptor traps near the channel/buffer interface followed by redistribution of these holes near the buffer/nucleation interface results in increase in negatively ionized buffer acceptor traps near the channel during stressing [10]. This increased negative charge near the channel results in channel depletion, leading to dynamic  $R_{ON}$ .

Substrate biasing is an effective tool in determining the dominant trapping mechanism under a given bias condition [5]. A negative  $V_{Sub}$  would oppose the injection of electrons into the GaN buffer, thereby reducing electron trapping in the GaN buffer and related dynamic  $R_{ON}$ . On the other hand, it would enhance redistribution of the holes emitted from the buffer acceptor traps to the buffer/nucleation interface, thereby increasing dynamic  $R_{ON}$  [10]. Fig. 5 shows reduction in dynamic  $R_{ON}$  of the HEMTs with negative  $V_{Sub}$ . The figure thus indicates electron trapping in the C-doped GaN buffer to be the major contributor of dynamic  $R_{ON}$  in GaN HEMTs.

In addition to substrate biasing, another way to determine electron trapping in the GaN buffer is to observe signatures of interaction of electrons with the buffer traps. Taking advantage of the fact that the device is stressed in semi-ON state,



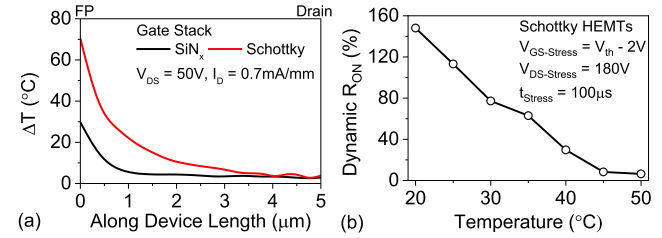


**Fig. 6.** (a)  $I_D$ - $V_{GS}$  of the devices measured under pristine condition and post a semi-ON state stressing ( $V_{GS-Stress} = V_{th} + 0.3$  V,  $V_{DS-Stress} = 180$  V,  $t_{Stress} = 1$  s). (b) EL line scans in gate-drain access region of the devices. (c) PL spectra of the GaN buffer obtained with 325 nm laser, showing the band-edge transitions, YL and BL bands. The EL spectra obtained at the FP edge of the SiN<sub>x</sub>- and Schottky-gated HEMTs show deviation (indicated by arrows) from the Maxwellian tail (indicated by dotted lines) in the YL and BL bands. (d) Dynamic  $R_{ON}$  of the Al<sub>0.5</sub>Ti<sub>0.5</sub>O<sub>y</sub>/SiN<sub>x</sub> passivated HEMTs with  $V_{DS-Stress}$ , as a function of  $t_{Stress}$ , under semi-ON state stressing.

EL spectrum analysis can be used to determine interaction of hot electrons with buffer traps [18]. EL line scans of SiN<sub>x</sub>- and Schottky-gated devices depict peak EL intensity near the FP edge, as shown in Fig. 6(b). This establishes EL hot spot to be near FP edge, where EL spectrum analysis should be carried out to determine interaction of hot electrons with buffer traps. Keeping above discussion in mind, EL spectrum of SiN<sub>x</sub>- and Schottky devices was captured near the FP edge during stressing the device in semi-ON state with a  $V_{DS-Stress}$  of 180 V and  $t_{Stress} > 30$  s. Such a stress condition was selected as it induced significant dynamic  $R_{ON}$  in both the devices. The results are as shown in Fig. 6(c). It shows that the EL intensity deviates from the Maxwellian tail-like features as we move toward higher energy values in the EL spectrum. This deviation is attributed to emergence of distinct peaks in the spectrum, indicated by arrows in the figure. A comparison of the EL spectra with the photoluminescence (PL) spectra of the GaN buffer indicates that these peaks correspond to the C-doping-induced yellow luminescence (YL) and blue luminescence (BL) bands [20], as seen in Fig. 6(c). The figure thus establishes interaction of hot electrons with the C-doping-induced buffer traps for both the SiN<sub>x</sub>- and Schottky-gated HEMTs. The above experiments thus establish trapping of electrons in the GaN buffer traps to be the major source of  $V_{cr}$  and high dynamic  $R_{ON}$  observed in these devices.

### C. Role of Electric Field

Discussion in the previous section established an interaction of hot electrons with buffer traps near the FP edge. To determine how this interaction affects dynamic  $R_{ON}$  behavior of the device, a relaxation in the EL intensity near the FP edge is required. To achieve this objective, devices with AlTiO surface protection over the SiN<sub>x</sub> passivation [4], as shown in Fig. 1(b), were fabricated. Surface charge modulation due to p-type nature of AlTiO surface protection effectively relaxes EL intensity near the FP edge [4], [21], as can be seen from EL intensity plots for AlTiO/SiN surface passivated devices in Fig. 6(b). The figure shows a uniform distribution of EL intensity in the entire gate-drain access region for devices with AlTiO/SiN surface passivation. Dynamic  $R_{ON}$  of these devices is shown in Fig. 6(d). The figure shows an absence of  $V_{cr}$  like phenomenon for these devices even for a  $t_{Stress}$  of 30 s. This establishes mitigation of electron trapping in the C-doped GaN buffer by EL intensity relaxation near the FP edge. As EL intensity is a direct measure of the magnitude of electric



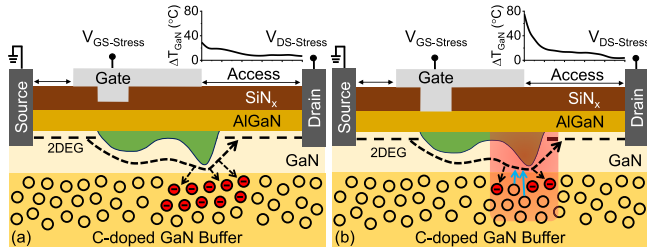
**Fig. 7.** (a) Average temperature ( $\Delta T$ ) profile of the GaN channel in the gate-drain access region of the devices, obtained using a thermo-reflectance setup. The  $\Delta T$  profile was measured as increase in device temperature as compared to ambient, under semi-ON state stressing with similar dissipated power. (b) Impact of temperature on OFF-state dynamic  $R_{ON}$  of Schottky HEMTs.

field, this further establishes that electric field relaxation near FP edge can mitigate electron trapping and hence  $V_{cr}$  in the dynamic  $R_{ON}$  behavior of AlGaIn/GaN HEMTs. The observed increase in dynamic  $R_{ON}$  values without a  $V_{cr}$  for higher values of  $V_{DS-Stress}$  and  $t_{Stress}$ , as seen in Fig. 6(d), can be attributed to hole emission from buffer traps [10], or trapping of hot electrons in the surface traps [12], [13].

### D. Role of Lattice Temperature

Discussion till now has established the role of electron trapping in GaN buffer and its dependence on channel electric field in determining the dynamic  $R_{ON}$  behavior of the devices. These parameters did not show any dependence on gate-stack, thereby suggesting a gate-stack independent dynamic  $R_{ON}$  behavior. However, Fig. 4 depicts a significantly different dynamic  $R_{ON}$  behavior of the Schottky- and SiN<sub>x</sub>-gated HEMTs. Explanation of this dependence thus needs an understanding of the impact of hot electrons on device parameters.

Apart from trapping processes, ionization of buffer acceptor traps can be affected by temperature distribution within the device. Stressing in semi-ON state can result in significant self-heating of the device due to increase in phonon population triggered by hot electron-lattice interaction [18]. This self-heating can modulate buffer trap ionization and hence will affect dynamic  $R_{ON}$  of the device. The channel temperature profile for the SiN<sub>x</sub>- and Schottky-gated HEMTs along the gate-drain access region is measured using a thermo-reflectance setup with 365 nm noncoherent LED light source and shown in Fig. 7(a). The devices were biased in semi-ON state with similar drain current for  $\sim 5$  s while



**Fig. 8.** Schematic showing the phenomenon governing semi-ON state stressing-induced dynamic  $R_{ON}$  in (a)  $\text{SiN}_x$ -gated and (b) Schottky HEMTs. The depletion region (electric field) in the devices near the FP edge results in injection of electrons into the GaN buffer of the HEMTs (shown by black arrows). Higher lattice temperature (shown by red window) of the Schottky HEMTs near the FP edge however results in de-trapping of carriers from buffer traps (shown by the blue arrows). This enhanced de-trapping results in lower net trapping in the GaN buffer of the Schottky devices.

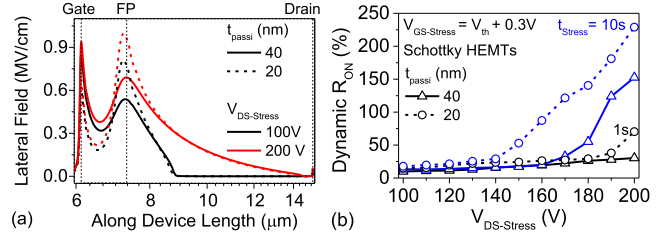
extracting the temperature profile. The figure shows a  $\sim 2.5\times$  higher temperature near the FP edge of the Schottky devices, as compared to the  $\text{SiN}_x$ -gated HEMTs. Higher temperature in Schottky HEMTs can be attributed to a stronger hot electron–GaN buffer interaction, as observed from Fig. 6(c). However, the phenomenon governing the observed higher temperature in Schottky HEMTs needs further investigation. Furthermore, higher temperature is known to increase hole emission [10] and should have increased dynamic  $R_{ON}$ . However, Fig. 4 shows a lower dynamic  $R_{ON}$  in Schottky HEMTs, as compared to  $\text{SiN}_x$ -gated HEMTs. This further establishes dynamic  $R_{ON}$  behavior to be not determined by hole emission.

To evaluate how temperature affects electron trapping-induced dynamic  $R_{ON}$  in Schottky devices, its dynamic  $R_{ON}$  behavior was tested at different temperatures in OFF-state. OFF-state stressing was selected here as it would induce minimum self-heating in the device due to lower currents, and also exhibits electron trapping-induced dynamic  $R_{ON}$  [5]. This allows determination of impact of temperature on electron injection and dynamic  $R_{ON}$ . Fig. 7(b) shows significant reduction in dynamic  $R_{ON}$  of the Schottky HEMTs under OFF-state stress, as device temperature is increased from 20 °C to 50 °C. Moreover, the HEMTs showed negligible dynamic  $R_{ON}$  when the device temperature was increased beyond 45 °C. This indicates a reduction in electron trapping-induced dynamic  $R_{ON}$  at higher temperature, which can be explained by thermal excitation-induced de-trapping of electrons trapped in the buffer acceptor traps.

The significantly high temperature near the FP edge of the Schottky HEMTs during semi-ON state stressing [see Fig. 7(a)] would hence result in similar de-trapping process. This reduces the dynamic  $R_{ON}$  of the devices as compared to the  $\text{SiN}_x$ -gated HEMTs, as seen in Fig. 4. The self-heating-induced de-trapping of electrons from buffer traps thus explains the gate-stack dependence of dynamic  $R_{ON}$  behavior of the devices.

## V. PHENOMENON GOVERNING DYNAMIC $R_{ON}$ UNDER SEMI-ON STATE STRESSING

Based on the above discussions, the phenomenon governing the semi-ON state stressing induced  $V_{cr}$  for dynamic  $R_{ON}$  in AlGaIn/GaN HEMTs is depicted in Fig. 8. The high electric



**Fig. 9.** (a) Lateral electric field profile in Schottky gated HEMTs, extracted using TCAD simulations, as a function of  $V_{DS}$ -Stress and  $\text{SiN}_x$  passivation thickness ( $t_{passi}$ ). (b) Experimentally extracted dynamic  $R_{ON}$  of Schottky HEMTs with varying  $t_{passi}$ , as a function of  $V_{DS}$ -Stress and  $t_{Stress}$ .

field near the FP edge of both the  $\text{SiN}_x$ -gated and Schottky HEMTs results in injection of electrons into the C-doped GaN buffer, as shown in Fig. 8(a) and (b) respectively. These electrons get trapped in the buffer traps and result in dynamic  $R_{ON}$  in the  $\text{SiN}_x$ -gated devices. However, the hot electron-induced self-heating in the channel of the Schottky HEMTs near the FP edge (where maximum trapping happens) results in increased electron de-trapping from these traps, as shown in Fig. 8(b). This enhanced de-trapping reduces the net electrons trapped in the GaN buffer and thus the dynamic  $R_{ON}$  of the Schottky HEMTs, as seen in Fig. 4.

Further, the  $t_{Stress}$  dependence of the dynamic  $R_{ON}$  behavior of the devices (see Fig. 4) can be explained as follows. An increase in dynamic  $R_{ON}$  requires substantial trapping in the GaN buffer [5]. This trapping would be determined by the number of electrons available in the buffer, probability of trapping and time available for trapping. While the electrons available in the buffer and probability of trapping is controlled by the electric field near FP edge or  $V_{DS}$ -Stress,  $t_{Stress}$  determines the time available for trapping. Increasing  $t_{Stress}$  thus relaxes the requirement of  $V_{DS}$ -Stress needed to obtain the same concentration of ionized buffer traps necessary for exhibiting increased dynamic  $R_{ON}$  in these devices. Hence, increasing  $t_{Stress}$  reduces the  $V_{cr}$  required for high dynamic  $R_{ON}$  in the  $\text{SiN}_x$ -gated HEMTs, as seen in Fig. 4(a). Moreover, enhanced de-trapping near the FP edge of the Schottky HEMTs, due to higher hot electron-induced self-heating, reduces the net electron trapping in the C-doped GaN buffer. Thus, a higher  $t_{Stress}$  and  $V_{DS}$ -Stress is required to achieve a similar concentration of ionized buffer traps so as to exhibit high dynamic  $R_{ON}$ , as seen in Fig. 4(b). Furthermore, increasing  $t_{Stress}$  reduces net traps available in the GaN buffer for electron trapping. This reduces electron trapping and can be balanced by Poole-Frenkel emission (PFE) from the ionized buffer traps at longer  $t_{Stress}$  ( $\sim$ few seconds) [5]. This balance of electron trapping and PFE results in the saturation of  $V_{cr}$  for  $t_{Stress}$  beyond  $\sim 10$ 's of seconds, as seen in Fig. 4.

The dependence of dynamic  $R_{ON}$  on electric field strength at FP edge ( $V_{DS}$ -Stress) and  $t_{Stress}$  can be validated by comparing dynamic  $R_{ON}$  behavior of the devices with different field magnitude near the FP edge. Computational estimation of electric field of devices with different  $\text{SiN}_x$  passivation thickness ( $t_{passi}$ ), as shown in Fig. 9(a), shows that this objective can be achieved by varying  $t_{passi}$ . Decreasing  $t_{passi}$  increases the electric field strength at FP edge, as seen in Fig. 9(a).

This increased electric field strength would lead to an increase in electron injection and subsequent trapping, thereby resulting in an increase in dynamic  $R_{ON}$  of the device with reduction in  $t_{passi}$ , as experimentally observed in Fig. 9(b). Moreover, Fig. 9(b) also shows a reduction in  $V_{cr}$  with reduction in  $t_{passi}$ . This dependence of dynamic  $R_{ON}$  and  $V_{cr}$  on field magnitude near FP edge validates the phenomenon proposed in Fig. 8.

## VI. CONCLUSION

Experimentations on Schottky and  $\text{SiN}_x$ -gated AlGaIn/GaN HEMTs on C-doped GaN buffer revealed a gate-stack-dependent dynamic  $R_{ON}$  behavior under semi-ON state stressing. Experiments showed a critical drain to source stress voltage beyond which dynamic  $R_{ON}$  increased significantly. However, it was observed at higher stress voltages and stress times for the Schottky HEMTs as compared to the  $\text{SiN}_x$ -gated devices. A similar critical voltage and dynamic  $R_{ON}$  behavior of both the devices under OFF-state stressing established the gate-stack dependence to be unique to semi-ON state stressing. Experiments showing reduction in dynamic  $R_{ON}$  with negative substrate bias established electron trapping in the C-doped GaN buffer to govern dynamic  $R_{ON}$  in both the devices. Moreover, interaction of hot electrons with buffer traps in EL spectra near the FP edge in both the devices established trapping to be happening near the FP edge. Thermoreflectance analysis however showed a significantly higher channel temperature near the FP edge of the Schottky HEMTs. This was attributed to self-heating induced by interaction of highly energetic hot electrons with the GaN buffer. The observed higher temperature was proposed to accelerate de-trapping from the C-doped GaN buffer and thereby reduce dynamic  $R_{ON}$ . Reduction in the dynamic  $R_{ON}$  of Schottky HEMTs under OFF-state stress (i.e., absence of channel hot electrons) at higher temperatures validated this proposal. Identical trapping due to similar electric field conditions but increased hot electron-induced self-heating and resulting de-trapping leads to a reduction in net ionized traps in the C-doped GaN buffer of the Schottky HEMTs. This results in requirement of a higher stress voltage and stress time for inducing high dynamic  $R_{ON}$  in the Schottky HEMTs, as compared to the  $\text{SiN}_x$ -gated AlGaIn/GaN HEMTs, under semi-ON-state stressing conditions.

## ACKNOWLEDGMENT

The authors would like to thank the National Nanofabrication Center (NNfC), Indian Institute of Science at Bangalore, Bengaluru, India, funded by MeitY and DST, Government of India. Sayak Dutta Gupta and Rajarshi Roy Chaudhuri would also like to thank DST INSPIRE for their fellowship.

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