

INDIAN INSTITUTE OF SCIENCE
E9-252: MATHEMATICAL METHODS AND TECHNIQUES IN SIGNAL PROCESSING
HOME WORK #3 - SOLUTIONS, FALL 2015

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Problem 1. If \mathcal{V} and \mathcal{W} are finite dimensional orthogonal subspaces of an inner product space \mathcal{H} , prove that $\dim(\mathcal{V} \oplus \mathcal{W}) = \dim(\mathcal{V}) + \dim(\mathcal{W})$.

Solution. Let $\dim(\mathcal{V}) = d_v$ and $\dim(\mathcal{W}) = d_w$. There exists a set of d_v vectors $\{\underline{v}_1, \underline{v}_2, \dots, \underline{v}_{d_v}\}$ that form an orthogonal basis for \mathcal{V} . Similarly, we can find a set of d_w vectors $\{\underline{w}_1, \underline{w}_2, \dots, \underline{w}_{d_w}\}$ that form an orthogonal basis for \mathcal{W} . Since the \mathcal{V} and \mathcal{W} are orthogonal subspaces, the orthogonal basis vectors to $\{\underline{v}_1, \underline{v}_2, \dots, \underline{v}_{d_v}, \underline{w}_1, \underline{w}_2, \dots, \underline{w}_{d_w}\}$ are all orthogonal to each other. Therefore, they all are linearly independent.

Any vector in $\mathcal{V} \oplus \mathcal{W}$ can be represented as a linear combination of two vectors \underline{v} and \underline{w} where $\underline{v} \in \mathcal{V}$ and $\underline{w} \in \mathcal{W}$. Since any vector in each of the subspaces can be represented as a linear combination of orthogonal basis vectors, any vector in $\mathcal{V} \oplus \mathcal{W}$ can be represented as a linear combination of vectors in $\{\underline{v}_1, \underline{v}_2, \dots, \underline{v}_{d_v}, \underline{w}_1, \underline{w}_2, \dots, \underline{w}_{d_w}\}$. Therefore, the set of vectors $\{\underline{v}_1, \underline{v}_2, \dots, \underline{v}_{d_v}, \underline{w}_1, \underline{w}_2, \dots, \underline{w}_{d_w}\}$ form an orthogonal basis for $\mathcal{V} \oplus \mathcal{W}$.

Therefore,

$$\dim(\mathcal{V} \oplus \mathcal{W}) = d_v + d_w = \dim(\mathcal{V}) + \dim(\mathcal{W}).$$

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Problem 2. Obtain the Haar wavelet decomposition of the signal $f(t)$. Indicate the signal dimension at each subspace carefully. Devise a generic algorithm for doing Haar decomposition using a computer program.

$$f(t) = \begin{cases} 2 & -2 \leq t < -1 \\ -4 & -1 \leq t < -0.5 \\ -2 & -0.5 \leq t < 0 \\ 2 & 0 \leq t < 0.25 \\ 1 & 0.25 \leq t \leq 2 \end{cases}$$

Solution. The Haar wavelet decomposition of the signal is $\underline{a}^{(0)} = (2, -3, 1.25, 1)$, $\underline{b}_0 = (0, -1, 0.25, 0)$, $\underline{b}_1 = (0, 0, 0, 0, 0.5, 0, 0, 0)$.

Dimension of the signal in \mathcal{V}_0 is 4, in \mathcal{V}_1 is 8 and in \mathcal{V}_2 is 16. The coefficients of Haar wavelet and scaling function in different subspaces are given below.

$a_k^{(j)}$	$k = -8$	-7	-6	-5	-4	-3	-2	-1	0	1	2	3	4	5	6	7
$j = 2$	2	2	2	2	-4	-4	-2	-2	2	1	1	1	1	1	1	1
$j = 1$					2	-4	-2	-2	1.5	1	1	1				
$j = 0$							2	-3	1.25	1						

$b_k^{(j)}$	$k = -4$	-3	-2	-1	0	1	2	3
$j = 1$	0	0	0	0	0.5	0	0	0
$j = 0$			0	-1	0.25	0		

TABLE 1. Coefficients of Haar wavelets and scaling function in different subspaces.

Figure 1 shows the decomposition of $f(t)$ as $v_0(t) + w_0(t) + w_1(t)$.

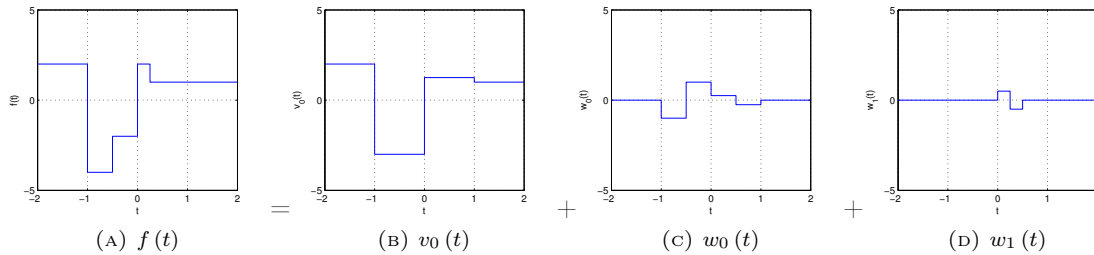


FIGURE 1. Haar decomposition of $f(t)$ as $v_0(t) + w_0(t) + w_1(t)$.

MATLAB code for wavelet decomposition is given below. The code uses the following relations that were derived in the class:

$$a_k^{(j-1)} = \frac{a_{2k}^{(j)} + a_{2k+1}^{(j)}}{2}$$

$$b_k^{(j-1)} = \frac{a_{2k}^{(j)} - a_{2k+1}^{(j)}}{2}$$

```

1 %% Signal Representation and Properties %%
2 intervals = [ -2, -1, -0.5, 0, 0.25, 2];
3 vals      = [ 2, -4, -2, 2, 1];
4
5 num_intervals      = length(intervals) - 1;
6 resolution         = 0.25;
7 duration           = intervals(end) - intervals(1);
8 num_decompositions = -log2(resolution);
9 signal_dim         = duration/resolution;
10
11 space_resolution   = (0.5).^(0:num_decompositions);
12 a = zeros(num_decompositions + 1, signal_dim);
13 b = zeros(num_decompositions, signal_dim);

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14
15 %% Signal representation using Haar scaling function %%
16 j = 1;
17 for i=1:(signal_dim)
18     end_time = intervals(1) + resolution*i;
19     if end_time <= intervals(j+1)
20         a(1, i) = vals(j);
21     else
22         j = j + 1;
23         a(1, i) = vals(j);
24     end
25 end
26
27 a(1,end) = vals(end);
28
29 %% Haar Wavelet Decomposition %%
30 for i=1:(num_decompositions)
31     a(i+1,1:(signal_dim/2)) = (a(i,1:2:signal_dim) + a(i, 2:2:signal_dim))/2;
32     b(i,1:(signal_dim/2)) = (a(i,1:2:signal_dim) - a(i, 2:2:signal_dim))/2;
33 end
34
35
36 %% Plotting the projections %%
37 space_resolution = resolution;
38 space_dim = signal_dim;
39 for i=1:(num_decompositions+1)
40     X = zeros(1,space_dim*2);
41     Y = zeros(1,space_dim*2);
42     X(1:2:end) = intervals(1) + (0:(space_dim-1))*space_resolution;
43     X(2:2:end) = intervals(1) + (1:space_dim)*space_resolution;
44     Y(1:2:end) = a(i, 1:space_dim);
45     Y(2:2:end) = a(i, 1:space_dim);
46     plot(X,Y); grid on; ylim([-5,5]); figure;
47     space_resolution = 2*space_resolution;
48     space_dim = space_dim/2;
49 end
50
51
52 %% Plotting the wavelet decompositions %%
53 space_resolution = resolution;
54 space_dim = signal_dim/2;
55 for i=1:(num_decompositions)
56     X = zeros(1,space_dim*4);
57     Y = zeros(1,space_dim*4);
58     X(1:2:end) = intervals(1) + (0:(2*space_dim-1))*space_resolution;
59     X(2:2:end) = intervals(1) + (1:(2*space_dim))*space_resolution;
60     Y(1:4:end) = b(i, 1:space_dim);
61     Y(2:4:end) = b(i, 1:space_dim);
62     Y(3:4:end) = -b(i, 1:space_dim);
63     Y(4:4:end) = -b(i, 1:space_dim);
64     plot(X,Y); grid on; ylim([-5,5]); figure;
65     space_resolution = 2*space_resolution;
66     space_dim = space_dim/2;
67 end

```



Problem 3. Prove the following properties for Haar wavelets:

- Parseval's equality i.e., energy conservation relation.
- Orthogonality across scales and time translates.

Solution. Orthogonality across scales and time translates

We have the Haar wavelet

$$\psi(t) = \begin{cases} 1, & 0 \leq t < \frac{1}{2} \\ -1, & \frac{1}{2} \leq t < 1 \\ 0 & \text{otherwise.} \end{cases}$$

The scaled and shifted versions are

$$\psi(2^j t - k) = \begin{cases} 1, & k2^{-j} \leq t < (k + \frac{1}{2})2^{-j} \\ -1, & (k + \frac{1}{2})2^{-j} \leq t < (k + 1)2^{-j} \\ 0 & \text{otherwise.} \end{cases}$$

Shift orthogonality: For $k \neq l$, $[k2^{-j}, (k + 1)2^{-j})$ and $[l2^{-j}, (l + 1)2^{-j})$ are non-overlapping regions. Therefore,

$$\langle \psi(2^j t - k), \psi(2^j t - l) \rangle = 0. \quad k \neq l$$

$$\langle \psi(2^j t - k), \psi(2^j t - k) \rangle = \int_{k2^{-j}}^{(k+1)2^{-j}} 1 dt = 2^{-j}.$$

Orthogonality across scales: Without loss of generality let, $p > q$ be two different scales. Therefore, $\psi(2^p t - k)$ has a smaller support than $\psi(2^q t - l)$ i.e., $2^{-p} < 2^{-q}$. Notice

1) $\psi(2^q t - l)$ is constant in each of the intervals $(m2^{-(q+1)}, (m + 1)2^{-(q+1)})$, $m \in \mathbb{Z}$.

2) For $p > q$, any interval $(n2^{-p}, (n + 1)2^{-p})$, $n \in \mathbb{Z}$ is a proper subset of the interval $(m2^{-(q+1)}, (m + 1)2^{-(q+1)})$ where $m = \lfloor n2^{q+1-p} \rfloor$.

Therefore,

$$\begin{aligned} \int_{k2^{-p}}^{(k+\frac{1}{2})2^{-p}} \psi(2^p t - k) \psi(2^q t - l) dt &= - \int_{(k+\frac{1}{2})2^{-p}}^{(k+1)2^{-p}} \psi(2^p t - k) \psi(2^q t - l) dt \\ \implies \langle \psi(2^p t - k), \psi(2^q t - l) \rangle &= \int_{k2^{-p}}^{(k+\frac{1}{2})2^{-p}} \psi(2^q t - l) dt - \int_{(k+\frac{1}{2})2^{-p}}^{(k+1)2^{-p}} \psi(2^q t - l) dt \\ &= 0 \end{aligned}$$

Therefore, the scales and time translates of Haar wavelets $\{\psi(2^j - k), k \in \mathbb{Z}, j = 0, 1, \dots\}$ are all orthogonal to each other.

Parseval's equality

Let $f(t) \in L^1(\mathbb{R})$. From the class notes, the wavelet decomposition of the function is given by

$$f(t) = v_0(t) + \sum_{j=0}^{\infty} w_j(t), \quad (1)$$

where

$$\begin{aligned} v_0(t) &= \sum_{k=-\infty}^{\infty} a_k^{(0)} \phi(t - k) \in \mathcal{V}_0 \\ v_j(t) &= \sum_{k=-\infty}^{\infty} b_k^{(j)} 2^{j/2} \psi(2^j t - k) \in \mathcal{W}_j \end{aligned}$$

and the coefficients $a_k^{(0)}$ and $b_k^{(j)}$ are obtained by projecting $f(t)$ onto the orthonormal basis of $\mathcal{V}_0, \mathcal{W}_0, \mathcal{W}_1, \mathcal{W}_2 \dots$, i.e.,

$$\begin{aligned} a_k^{(0)} &= \langle f(t), \phi(t - k) \rangle, \\ b_k^{(k)} &= \langle f(t), 2^{j/2} \psi(2^j t - k) \rangle. \end{aligned}$$

Our aim is to prove that the energy in the signal is equal to the sum of square of the wavelet coefficients $\{a_k^{(0)}, b_k^{(j)}, k \in \mathbb{Z}, j = 0, 1, \dots\}$ i.e.,

$$\int_{-\infty}^{\infty} |f(t)|^2 dt = \sum_{k=-\infty}^{\infty} |a_k^{(0)}|^2 + \sum_{j=0}^{\infty} \sum_{k=-\infty}^{\infty} |b_k^{(j)}|^2$$

The subspaces $\mathcal{V}_0, \mathcal{W}_0, \mathcal{W}_1, \dots$ are all orthogonal to each other. Therefore,

$$\begin{aligned} \langle v_0(t), w_j(t) \rangle &= 0, \quad j = 0, 1, 2, \dots \\ \langle w_j(t), w_k(t) \rangle &= 0, \quad j \neq k. \end{aligned}$$

Therefore,

$$\begin{aligned} \langle v_0(t), f(t) \rangle &= \langle v_0(t), v_0(t) \rangle + \sum_{j=0}^{\infty} \langle v_0(t), w_j(t) \rangle = \langle v_0(t), v_0(t) \rangle \\ \langle w_j(t), f(t) \rangle &= \langle w_j(t), v_0(t) \rangle + \sum_{j=0}^{\infty} \langle w_j(t), w_j(t) \rangle = \langle w_j(t), w_j(t) \rangle \end{aligned}$$

Hence,

$$\int_{-\infty}^{\infty} |f(t)|^2 dt = \langle f(t), f(t) \rangle = \langle v_0(t), v_0(t) \rangle + \sum_{j=0}^{\infty} \langle w_j(t), w_j(t) \rangle. \quad (2)$$

Consider

$$\begin{aligned} \langle v_0(t), v_0(t) \rangle &= \left\langle \sum_{k=-\infty}^{\infty} a_k^{(0)} \phi(t-k), \sum_{l=-\infty}^{\infty} a_l^{(0)} \phi(t-l) \right\rangle \\ &= \sum_{k=-\infty}^{\infty} \sum_{l=-\infty}^{\infty} a_k^{(0)} \overline{a_l^{(0)}} \langle \phi(t-k), \phi(t-l) \rangle \\ &= \sum_{k=-\infty}^{\infty} \sum_{l=-\infty}^{\infty} a_k^{(0)} \overline{a_l^{(0)}} \delta_{kl} = \sum_{k=-\infty}^{\infty} |a_k^{(0)}|^2. \end{aligned} \quad (3)$$

Similarly,

$$\begin{aligned} \langle w_j(t), w_j(t) \rangle &= \left\langle \sum_{k=-\infty}^{\infty} b_k^{(j)} 2^{j/2} \psi(2^j t - k), \sum_{l=-\infty}^{\infty} b_l^{(j)} 2^{j/2} \psi(2^j t - l) \right\rangle \\ &= \sum_{k=-\infty}^{\infty} \sum_{l=-\infty}^{\infty} b_k^{(j)} \overline{b_l^{(j)}} \langle 2^{j/2} \psi(2^j t - k), 2^{j/2} \psi(2^j t - l) \rangle \\ &= \sum_{k=-\infty}^{\infty} \sum_{l=-\infty}^{\infty} b_k^{(j)} \overline{b_l^{(j)}} \delta_{kl} = \sum_{k=-\infty}^{\infty} |b_k^{(j)}|^2. \end{aligned} \quad (4)$$

From (2), (3) and (4), we have the Parseval's equality

$$\int_{-\infty}^{\infty} |f(t)|^2 dt = \sum_{k=-\infty}^{\infty} |a_k^{(0)}|^2 + \sum_{j=0}^{\infty} \sum_{k=-\infty}^{\infty} |b_k^{(j)}|^2.$$

■

Remark:

- Notice that the Parseval's equality is same as the following result that we have proved for inner product spaces: If $\{\underline{v}_1, \underline{v}_2, \dots\}$ form orthonormal basis for an inner product space \mathcal{V} and a vector $\underline{v} \in \mathcal{V}$ is written as $\underline{v} = \sum_{k=1}^{\infty} a_k \underline{v}_k$, then

$$|\underline{v}|^2 = \langle \underline{v}, \underline{v} \rangle = \sum_{k=1}^{\infty} |a_k|^2.$$

Problem 4. For $j \in \mathbb{Z}$, let \mathcal{V}_j be the space of all signals $f(t) \in L^2$ bandlimited within the interval $[-2^j\pi, 2^j\pi]$. Consider the signal $\phi(t) := \text{sinc}(t) = \frac{\sin(\pi t)}{\pi t}$. Prove the following.

- The nesting, closure, shrinking and scaling properties that we discussed in the class as part of the multiresolution analysis definition.
- $\{\phi(t - k), k \in \mathbb{Z}\}$ is a shift orthogonal basis for \mathcal{V}_0 .
- $\phi(t) = \phi(2t) + \sum_{k \in \mathbb{Z}} \frac{2(-1)^k}{(2k+1)\pi} \phi(2t - 2k - 1)$. (Scaling relation)

Solution. Part 1

Nesting property: If $f(t) \in \mathcal{V}_j$, then $f(t)$ is bandlimited within the interval $[-2^j\pi, 2^j\pi] \implies f(t)$ is bandlimited within the interval $[-2^{j+1}\pi, 2^{j+1}\pi] \implies f(t) \in \mathcal{V}_{j+1}$. Therefore, $\mathcal{V}_j \subset \mathcal{V}_{j+1}$, $j = 0, 1, \dots$. Therefore

$$\mathcal{V}_0 \subset \mathcal{V}_1 \subset \mathcal{V}_2 \subset \mathcal{V}_3 \dots$$

Scaling property: If $f(t) \in \mathcal{V}_j$, then $f(t)$ is bandlimited within the interval $[-2^j\pi, 2^j\pi] \implies f(2t)$ is bandlimited within the interval $[-2^{j+1}\pi, 2^{j+1}\pi] \implies f(2t) \in \mathcal{V}_{j+1}$. Therefore, $\{f(2t) \mid f(t) \in \mathcal{V}_j\} \subseteq \mathcal{V}_{j+1}$.

Similarly, if $f(t) \in \mathcal{V}_{j+1}$, then $f(t)$ is bandlimited within the interval $[-2^{j+1}\pi, 2^{j+1}\pi] \implies f(\frac{t}{2})$ is bandlimited within the interval $[-2^j\pi, 2^j\pi] \implies f(\frac{t}{2}) \in \mathcal{V}_j$. Therefore, $\mathcal{V}_j \supset \{f(\frac{t}{2}) \mid f(t) \in \mathcal{V}_{j+1}\} \implies \{f(2t) \mid f(t) \in \mathcal{V}_j\} \supseteq \mathcal{V}_{j+1}$.

Therefore, we have $\{f(2t) \mid f(t) \in \mathcal{V}_j\} = \mathcal{V}_{j+1}$.

Shrinking property: Signals in \mathcal{V}_{-j} are bandlimited within the interval $[-2^{-j}\pi, 2^{-j}\pi]$ for $j \geq 0$. Therefore, signals in $\bigcap_{j \in \mathbb{Z}} \mathcal{V}_j$ will have only a d.c. component in the signal. The only square integrable d.c. signal is $f(t) = 0$.

Therefore, $\bigcap_{j \in \mathbb{Z}} \mathcal{V}_j = \{0\}$.

Closure property: Let $f(t) \in L^2(\mathbb{R})$ be any function in L^2 space. Consider its Fourier transform $F(\omega)$. Let $f_j(t)$ be the projection of $f(t)$ onto \mathcal{V}_j . Then its Fourier transform is

$$F_j(\omega) = \begin{cases} F(\omega) & \omega \in [-2^{-j}\pi, 2^{-j}\pi] \\ 0 & \text{otherwise.} \end{cases}$$

For the closure property, we need to show that $\lim_{j \rightarrow \infty} f_j(t) = f(t)$ in L^2 sense.

Split the frequencies into disjoint intervals given by $I_j = [-2^{-(j+1)}\pi, 2^{-(j+1)}\pi] \setminus [-2^{-j}\pi, 2^{-j}\pi]$, $j = 1, 2, \dots$ and $I_0 = [-\pi, \pi]$. We have $\bigcup_{j=0}^N I_j = [-2^{-(N+1)}\pi, 2^{-(N+1)}\pi]$ and $\bigcup_{j=0}^{\infty} I_j = \mathbb{R}$.

Let E be the energy in the signal $f(t)$. Using Parseval's theorem, we have

$$E = \int_{-\infty}^{\infty} |f(t)|^2 dt = \int_{-\infty}^{\infty} |F(\omega)|^2 d\omega = \sum_{j=0}^{\infty} \int_{\omega \in I_j} |F(\omega)|^2 d\omega. \quad (5)$$

Let E_j be the energy of the signal in the frequencies I_j i.e.,

$$E_j = \int_{\omega \in I_j} |F(\omega)|^2 d\omega, \quad j = 0, 1, 2, \dots$$

Then, the equation (5) can be written as

$$\begin{aligned} E &= \sum_{j=0}^{\infty} E_j \\ \implies \lim_{N \rightarrow \infty} \left(E - \sum_{j=0}^N E_j \right) &= 0. \quad (\because E \text{ is finite}) \end{aligned} \quad (6)$$

Notice that $\sum_{j=0}^N E_j$ is the energy of the signal $f(t)$ within the frequencies $[-2^{-(N+1)}\pi, 2^{-(N+1)}\pi]$. Therefore, $\sum_{j=0}^N E_j$ is the energy in $F_{N+1}(\omega)$. Therefore, $\left(E - \sum_{j=0}^N E_j \right)$ is the energy in $F(\omega) - F_{N+1}(\omega)$ ($= f(t) - f_{N+1}(t)$ in the time-domain). From equation (6),

$$\lim_{N \rightarrow \infty} \int_{-\infty}^{\infty} |f(t) - f_{N+1}(t)|^2 dt = 0$$

$$\implies \lim_{N \rightarrow \infty} f_N(t) = f(t) \quad (\text{in } L^2 \text{ sense})$$

Hence, any function in $L^2(\mathbb{R})$ can be represented within $\bigcup_{j=0}^N \mathcal{V}_j$. Therefore, $\bigcup_{j=0}^N \mathcal{V}_j = L^2(\mathbb{R})$.

Part 2:

We need to prove that

- a) $\{\phi(t-k), k \in \mathbb{Z}\}$ are orthogonal and
- b) $\{\phi(t-k), k \in \mathbb{Z}\}$ form basis for \mathcal{V}_0

To prove that $\{\phi(t-k), k \in \mathbb{Z}\}$ are orthogonal, consider the inner product between $\phi(t-k)$ and $\phi(t-l)$ for $k \neq l$,

$$\begin{aligned} \int_{-\infty}^{\infty} \phi(t-k) \phi(t-l) &= \int_{-\infty}^{\infty} \frac{\sin(\pi(t-k))}{\pi(t-k)} \frac{\sin(\pi(t-l))}{\pi(t-l)} dt \\ &= \int_{-\infty}^{\infty} (-1)^k \frac{\sin(\pi t)}{\pi(t-k)} (-1)^l \frac{\sin(\pi t)}{\pi(t-l)} dt \quad (\sin(t-n\pi) = (-1)^n \sin(t)) \\ &= \frac{(-1)^{l+k}}{\pi^2(k-l)} \int_{-\infty}^{\infty} \left(\frac{\sin^2(\pi t)}{t-k} - \frac{\sin^2(\pi t)}{t-l} \right) dt \\ &= \frac{(-1)^{l+k}}{\pi^2(k-l)} \int_{-\infty}^{\infty} \frac{\sin^2(\pi(t-k))}{t-k} dt \\ &\quad - \frac{(-1)^{l+k}}{\pi^2(k-l)} \int_{-\infty}^{\infty} \frac{\sin^2(\pi(t-l))}{t-l} dt \quad (\sin^2(t-n\pi) = \sin^2(t)) \\ &= \frac{(-1)^{l+k}}{\pi^2(k-l)} \int_{-\infty}^{\infty} \frac{\sin^2(\pi t)}{t} dt \quad (\text{Change of variables } t-k \rightarrow t) \\ &\quad - \frac{(-1)^{l+k}}{\pi^2(k-l)} \int_{-\infty}^{\infty} \frac{\sin^2(\pi t)}{t} dt \quad (\text{Change of variables } t-l \rightarrow t) \\ \int_{-\infty}^{\infty} \phi(t-k) \phi(t-l) &= 0 \end{aligned}$$

Proving that $\{\phi(t-k), k \in \mathbb{Z}\}$ form basis for \mathcal{V}_0 following from Nyquist sampling theorem. Any bandlimited signal can be sampled at Nyquist rate without losing any information. The original signal can be constructed from the samples using "sinc-interpolation".

Consider any signal $f(t) \in \mathcal{V}_0$ that is bandlimited to $[-\pi, \pi]$. The Nyquist rate is $2 \times \frac{\pi}{2\pi} = 1$ samp/sec. The sampled signal is

$$f_s(t) = \sum_{k \in \mathbb{Z}} f(k) \delta(t-k).$$

If $F(\omega)$ is the Fourier transform of $f(t)$, then the Fourier transform of $f_s(t)$ is same as $F(\omega)$ repeated at 2π intervals. Therefore, the signal $f(t)$ can be obtained from $f_s(t)$ by using an ideal low pass filter bandlimited to $[-\pi, \pi]$. In time-domain, this filter is given by $\phi(t) = \frac{\sin(\pi t)}{\pi t}$. Therefore,

$$\begin{aligned} f(t) &= \sum_{j=-\infty}^{\infty} f_s(t-j) \phi(j) \\ &= \sum_{j=-\infty}^{\infty} \sum_{k \in \mathbb{Z}} f(k) \delta(t-j-k) \phi(j) \\ &= \sum_{k \in \mathbb{Z}} f(k) \sum_{j=-\infty}^{\infty} \delta(t-k-j) \phi(j) \\ f(t) &= \sum_{k \in \mathbb{Z}} f(k) \phi(t-k) \quad (\phi(t) * \delta(t-k) = \phi(t-k)) \end{aligned}$$

Therefore, any function in \mathcal{V}_0 can be represented as a linear combination of $\{\phi(t-k), k \in \mathbb{Z}\}$. Hence, $\{\phi(t-k), k \in \mathbb{Z}\}$ form orthogonal basis for \mathcal{V}_0 .

Part 3

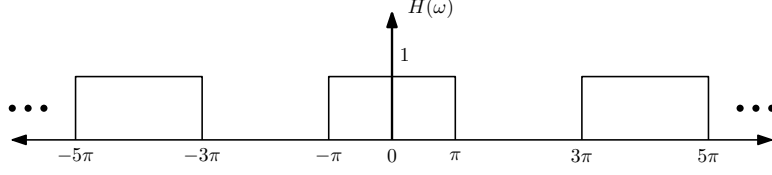


FIGURE 2. Filtering $\phi(2t)$ by this filter $H(\omega)$ results in $\phi(t)$

$\phi(t)$ is an ideal low pass filter band limited to $[-\pi, \pi]$. Its scale $\phi(2t)$ an ideal low pass filter band limited to $[-2\pi, 2\pi]$. $\phi(t)$ can be obtained from $\phi(2t)$ by passing through a filter $h(t)$ whose frequency response is as shown in Figure 2. The signal $h(t)$ is given by (derivation for this is given later)

$$h(t) = \delta(t) + \sum_{k \in \mathbb{Z}} \frac{2(-1)^k}{(2k+1)\pi} \delta(t - 2k - 1). \quad (7)$$

Therefore,

$$\begin{aligned} \phi(t) &= h(t) * \phi(2t) \\ &= \delta(t) * \phi(2t) + \sum_{k \in \mathbb{Z}} \frac{2(-1)^k}{(2k+1)\pi} \delta(t - 2k - 1) * \phi(2t) \\ \phi(t) &= \phi(2t) + \sum_{k \in \mathbb{Z}} \frac{2(-1)^k}{(2k+1)\pi} \phi(2t - 2k - 1) \end{aligned}$$

Proof for equation (7):

The frequency response in Figure 2, is an even periodic function with period 2π . Therefore, we can write the frequency response as

$$\begin{aligned} H(\omega) &= a_0 + \sum_{n=1}^{\infty} a_n \cos\left(n \frac{\omega}{2}\right), \\ \text{where, } a_0 &= \int_{-\pi}^{\pi} H(\omega) d\omega = 2\pi \\ a_n &= \int_{-\pi}^{\pi} H(\omega) \cos\left(n \frac{\omega}{2}\right) d\omega \\ &= \int_{-\pi}^{\pi} \cos\left(n \frac{\omega}{2}\right) d\omega \\ &= \frac{2}{n} \sin\left(n \frac{\omega}{2}\right) \Big|_{-\pi}^{\pi} \\ &= \frac{4}{n} \sin\left(n \frac{\pi}{2}\right) \\ &= \begin{cases} 0 & n \text{ is even} \\ (-1)^k \frac{4}{2k+1}, & n = 2k+1 \text{ is odd} \end{cases} \end{aligned}$$

Therefore,

$$\begin{aligned} H(\omega) &= 2\pi + \sum_{k=0}^{\infty} (-1)^k \frac{4}{2k+1} \cos\left((2k+1) \frac{\omega}{2}\right) \\ &= 2\pi + \sum_{k=0}^{\infty} (-1)^k \frac{2}{2k+1} \left(e^{j(2k+1)\frac{\omega}{2}} + e^{-j(2k+1)\frac{\omega}{2}} \right) \\ \implies h(t) &= \delta(t) + \sum_{k=0}^{\infty} (-1)^k \frac{1}{(2k+1)\pi} (\delta(t - 2k - 1) + \delta(t + 2k + 1)) \\ &= \delta(t) + \sum_{k=0}^{\infty} (-1)^k \frac{1}{(2k+1)\pi} \delta(t - 2k - 1) + \sum_{k=0}^{\infty} (-1)^k \frac{1}{(2k+1)\pi} \delta(t + 2k + 1) \end{aligned}$$

$$\begin{aligned}
&= \delta(t) + \sum_{k=0}^{\infty} (-1)^k \frac{1}{(2k+1)\pi} \delta(t-2k-1) \\
&\quad + \sum_{l=-\infty}^{-1} (-1)^{-1-l} \frac{1}{(2(-l-1)+1)\pi} \delta(t+2(-l-1)+1) \\
&= \delta(t) + \sum_{k=0}^{\infty} (-1)^k \frac{1}{(2k+1)\pi} \delta(t-2k-1) + \sum_{l=-\infty}^{-1} (-1)^{l+1} \frac{2}{-(2l+1)\pi} \delta(t-2l-1) \\
&= \delta(t) + \sum_{k=0}^{\infty} (-1)^k \frac{1}{(2k+1)\pi} \delta(t-2k-1) + \sum_{l=-\infty}^{-1} (-1)^l \frac{2}{(2l+1)\pi} \delta(t-2l-1) \\
&= \delta(t) + \sum_{k=-\infty}^{\infty} (-1)^k \frac{1}{(2k+1)\pi} \delta(t-2k-1)
\end{aligned}$$

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